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Reprinted from the JOURNAL OF CHEMICAL PHYSICS, Vol. 32, No. 5, 1418-1422, May, 1960 Printed in U. S. A.

# Effect of Pressure on the Spectra of Color Centers\*

EPPL RA60 0118

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#### (Received October 20, 1959)

The effect of pressure on the spectra of certain color centers produced in the alkali halides by x irradiation has been measured to pressures as high as 166 000 atm.

For the F center a shift to higher energies with increasing pressure is observed. The slope of the shift versus density is at least twice the value that would be predicted from Ivey's relation at low pressure, and decreases with increasing pressure. This indicates that the compressibility in the neighborhood of the F center is greater than in the bulk crystal, particularly at low pressure.

For the M center in LiCl a shift to higher energy with increasing pressure is observed, about one-fifth as great as the shift observed for the F center.

### INTRODUCTION

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THE effect of pressure on the spectra of certain color centers in alkali halide lattices produced by ionizing radiation has been studied to as high as 166 000 atm.

There have been three previous investigations of the effect of pressure on color centers in alkali halides. Burstein et al.<sup>1</sup> measured the shift of the F-band peak frequency to 2000 atm in three alkali halides. Jacobs<sup>2</sup> measured the effect of pressure on the peak frequency of the F band to as high as 8000 atm in seven alkali halides. Maisch and Drickamer<sup>3</sup> studied the effect of pressure on the spectra of the F band in six alkali halides and the M band in one other alkali halide to 50 000 atm.

They found that, in general, a shift to higher energy occurred for all bands with pressure as long as no phase transition occurred. Maisch and Drickamer<sup>3</sup> also found that in some alkali halides, there appears at higher pressures a new band on the high energy side of the Fband which they called the  $K^1$  band.

#### EXPERIMENTAL PROCEDURE

The types of crystals and color centers studied, the techniques used to introduce the centers, and the pressure ranges involved are summarized in Table I.

The crystals to be x irradiated were cleaved to about  $2\frac{1}{2} \times 5 \times \frac{1}{2}$  mm and exposed at a distance of one inch for times ranging from one to twenty-five hours to an x-ray tube with a tungsten target and beryllium windows operated at 40 kv and 15 ma. The crystals were stored in the dark until used. They were loaded into the bomb in a very dim red light. This procedure was found by Maisch and Drickamer<sup>3</sup> to give results identical to those obtained with additively colored crystals.

The crystals to be bleached were exposed from 5-10 min at a distance of 2-3 in. to the radiation from a

\* This work was supported in part by U. S. Atomic Energy Commission contract, Chemical Engineering Project 5. <sup>1</sup> E. Burstein, J. J. Oberly, and J. W. Davisson, Phys. Rev. 85, 729 (1952).

 <sup>2</sup> I. S. Jacobs, Phys. Rev. 93, 993 (1954).
<sup>3</sup> W. G. Maisch and H. G. Drickamer, J. Phys. Chem. Solids 5, 328 (1958).

Hanovia type 16 200 mercury lamp, then used immediately.

The high-pressure optical system and techniques have been described previously.4

# RESULTS AND DISCUSSION

# A. F Center in the fcc Structure

The primary result of x irradiation of a pure alkali halide crystal is the introduction of a band in the visible called the F center. The commonly accepted model of the defect, due to deBoer,<sup>5</sup> is that of an electron trapped at a negative ion vacancy. The shift of this peak with pressure is shown in Figs. 1 to 6 for six alkali halide crystals which have the fcc (NaCl) structure, at least at atmospheric pressure. The maximum scatter in the data is  $\pm 300$  cm<sup>-1</sup>. Figures 7 to 10 show typical spectra for KCl, KBr, and KI.6

It was first suggested by Burstein and Oberly<sup>7</sup> that the F center could, to the zeroth order, be treated as a particle trapped in a box. One of the important conclusions of this model is that  $E \propto (1/l)^2$  where l is the box diameter. Thus

# $\nu_m = E/h \propto (1/A_0^2) \alpha \rho^3$

where  $A_0$  is the lattice parameter and  $\epsilon$  is the density.

Mollwo,<sup>8</sup> using data at one atmosphere and 25°C for several alkali halides having the fcc structure, showed that  $\nu_m A_0^2$  is essentially constant. Later Ivey<sup>9</sup> correlated the data and found that

## $E_m A_0^{1.84} = 1.76 \times 10^{-19} \text{ v m}^2$ .

The density of certain alkali halides has been measured as a function of pressure to 100 000  $\rm kg/cm^2$  by

J. H. deBoer, Rec. trav. chim. 56, 301 (1937).

<sup>6</sup> It may be noted that the half-widths do not agree particularly well with other published values.

7 E. Burstein and J. J. Oberly, Natl. Bur. Standards Circ. No. 5109, 285 (1952).

<sup>8</sup> E. Mollwo, Nachr. Ges. Wiss. Göttingen, Math. phys. Kl. 97 (1931). <sup>9</sup> H. F. Ivey, Phys. Rev. 72, 341 (1947).

<sup>&</sup>lt;sup>4</sup> R. A. Fitch, T. E. Slykhouse, and H. G. Drickamer, J. Opt. Soc. Am. 47, 1015 (1957).

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FIG. 8. Color center spectra at various pressures-KBr.

the K' band in any case. This indicates that the formation of the K' band is related in some way to the transition between the two alkali halide structures.

(b) The paper by Lazarus<sup>11</sup> shows that the shear constant of sodium chloride is relatively high and is relatively independent of pressure, while the shear constant of potassium chloride is not only much lower at one atmosphere, but decreases with increasing pressure.

(c) Wiederhorn<sup>12</sup> has shown that the transmittance of the potassium halides decreases by a factor of over  $10^{-3}$  before the pressure induced phase transition is more than 10% complete.

The preceding indicate that the K' band growth may possibly by connected with the propagation of shear through the lattice, causing a disruption in the structure of the F center. It is possible that the defect ties down the ends of dislocations formed by the shear, introducing new forces into the energy scheme of the center.

# C. Transition Effects, and the F Center in the sc (CsCl) Structure

Three potassium halides (KCl, KBr, and KI) transform from the fcc to the sc structure at 19 000-



FIG. 9. Color center spectra at various pressures-KI.

<sup>12</sup> S, Wiederhorn (private communication).



Fig. 10.  $\log(\nu/\nu_0)_m \text{ vs } \log \rho / \rho_0 - \text{NaCl}; - - -, \nu/\nu_0 - (\rho/\rho)^{2/8}$ .



FIG. 11.  $\log(\nu/\nu_0)_m$  vs  $\log \rho/\rho_0$ —NaBr; — – — – ,  $\nu/\nu_0 \sim \rho/\rho_0)^{2/3}$ .



FIG. 12.  $\log(\nu/\nu_0)_m$  vs  $\log \rho/\rho_0 - KCl; - - -, \nu/\nu_0 - (\rho/\rho_0)^{2/3}$ .

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<sup>&</sup>lt;sup>11</sup> D. Lazarus, Phys. Rev. 76, 545 (1949).

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FIG. 13. Log(v/v0)m vs logp/p0-KBr; - $\sim (\rho/\rho_0)^{2/3}$ 



FIG. 14.  $\log(\nu/\nu_0)_m$  vs  $\log \rho/\rho_0$ —KI; — — —,  $\nu/\nu_0 \sim (\rho/\rho_0)^{2/3}$ .

20 000 atm pressure. A discontinuous shift of the peak frequency to higher energies is observed at the transition (see Figs. 3-5 and 12-14). Interpretation of these data are hampered by the growth of the K' band in KBr and KI, and by further structure in KI which also was observed by Maisch and Drickamer.3

# D. Effect of Pressure on the M Center

When crystals containing F centers are irradiated with light in the F-band region at room temperature, the F band bleaches and several new bands appear on the low energy side of the F band.<sup>13</sup> The strongest of these is the M center.

The nature of this center is uncertain. Seitz<sup>14</sup> proposed that it consists of an F center plus a vacancy pair. More recent experiments<sup>15</sup> indicate that it posses-

| TABLE II. Comparison of | f experimental and | I calculated values | of fr. |
|-------------------------|--------------------|---------------------|--------|
|-------------------------|--------------------|---------------------|--------|

| Compound | $f_T$ exptl.                                 | $f_T$ calc   |   |
|----------|--|--|---|
| NaCl     | 2.1  | 2.9  |   |
| NaBr     | 2.3  | 2.4  |   |
| KCl      | 2.3  | 2.1  |   |
| KBr      | 1.8  | 2.1  |   |
| KI       | 1.7  | 2.3  |   |
|          | Compound<br>NaCl<br>NaBr<br>KCl<br>KBr<br>KI | Compound $f_T$ exptl.       NaCl     2.1       NaBr     2.3       KCl     2.3       KBr     1.8       KI     1.7 | Compound $f_T$ exptl. $f_T$ calc       NaCl     2.1     2.9       NaBr     2.3     2.4       KCl     2.3     2.1       KBr     1.8     2.1       KI     1.7     2.3 |

TABLE III. Density of LiCl vs pressure.

| $P (kg/cm^2)$ |         | $\Delta V/V_0$ |  |
|---------------|---------|----------------|--|
|               | 0       | 0.000          |  |
|               | 50 000  | 0.121          |  |
|               | 100 000 | 0.183          |  |



FIG. 15. Effect of pressure on the spectra of the F and Mcenters in LiCl.

ses a higher degree of symmetry than Seitz' model would indicate. Knox16 proposed a modification of Seitz' model which possesses this higher symmetry, wherein a positive ion and an electron occupy a vacancy aggregation consisting of two negative ion vacancies and two positive ion vacancies.

The effect of pressure to 50 000 atm has been measured on the M center in LiCl. The results are compared with data for the F center in LiCl in Fig. 15. The shift with pressure for the M center is less than one-fifth that of the F center.

The low compressibility seems inconsistent with the Seitz<sup>14</sup> model. On the other hand, the Knox modification to this model<sup>16</sup> does appear to be somewhat more consistent with the data since the atom in the middle of the center might tend to oppose the compression to a greater degree than one fixed in the corner.

## ACKNOWLEDGMENTS

R. A. Eppler would like to thank the Shell Oil Company and the Dow Chemical Company for fellowship support.

<sup>16</sup> R. S. Knox, Phys. Rev. Letters 2, 87 (1959).

 <sup>&</sup>lt;sup>13</sup> F. Seitz, Revs. Modern Phys. 26, 7 (1954).
<sup>14</sup> F. Seitz, Revs. Modern Phys. 18, 384 (1946).
<sup>15</sup> A. W. Overhauser and H. Ruchardt, Phys. Rev. 112, 722 (1958).